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RESEARCH ARTICLE

VIBRATIONAL SPECTROSCOPIC STUDY AND NBO ANALYSIS ON TRANEXAMIC ACID USING DFT METHOD

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INTRODUCTION

Tranexmic acid (TA) is chemically 4-(aminomethyl) cyclohexane-1-carboxylic acid is a synthetic derivative of the amino acid lysine that exerts its antifibrinolytic effect through the reversible blockade of the lysine binding sites on plasminogen molecules (Thoresen et al., 1981). TA is useful in a wide range of haemorrhagic conditions. The drug reduces postoperative blood losses and transfusion requirements in a number types of surgery (Brown et al., 1997 and Ido et al., 2004), with potential cost and tolerability advantages over aprotinin, and appears to reduce rates of mortality and urgent surgery in patients with upper gastrointestinal haemorrhage. TA reduces menstrual blood loss and is a possible alternative to surgery in menorrhagia and has been used successfully to control bleeding in pregnancy. In the present study, FT-IR, FT-Raman spectral investigation of TA has been performed using density functional theory (DFT). The redistribution of electron density (ED) in various bonding and antibonding orbital and E(2) energies have been calculated by natural bond orbital (NBO) analysis by DFT method to give clear evidence of stabilization originating from the hyper conjugation of various intra-molecular interactions. The HOMO and LUMO analysis have been used to elucidate information regarding charge transfer within the molecule.

MATERIALS AND METHODS

The compound TA was purchased from sigma-Aldrich chemical company (USA) with a stated purity of greater than 97% and it was used as such without further purification. The FT-Raman spectrum of TA has been recorded using 1064 nm

ABSTRACT

In this work, we report a combined experimental and theoretical study on molecular structure, Vibrational spectra and NBO analysis of tranexamic acid (TA). The FT-Raman and FT-IR spectra of TA wear recorded in the solid phase. The molecular geometry, harmonic vibrational frequencies and bonding features of TA in the ground state have been calculated by using density functional method (B3LYP) with standard 6-31G (d,p) basis set. Stability of the molecule arising from hyper conjugative interactions, charge delocalization has been analyzed using natural bond orbital (NBO) analysis. The calculated HOMO and LUMO energies show that charge transfer occurs within the molecule. The theoretical FT-IR and FT-Raman spectra for the title molecule have been constructed.

line of Nd-YAG laser as excitation wavelength in the region 100-4000 cm⁻¹ on BrukerModel IFS 66 spectrophotometer. The FT-IR spectrum of this compound was recorded in the region 400-4000 cm⁻¹ on IFS 66V Spectrophotometer using KBr pellet technique with a scanning speed of 30 cm⁻¹ min⁻¹ and the spectral resolution of 4.0 cm⁻¹. The observed experimental and calculated FT-IR and FT-Raman spectra are shown in Figs. 1 and 2. The spectral measurements were carried out at Sophisticated Analytical Instrumentation Facility (SAIF), IIT, and Chennai.

Computational details

The entire calculations was performed at Density functional theory (DFT) levels with a Pentium Intel (R) core 2 quard 2.40 GHz personal computer using Gaussian 03W (Frisch et al.,2004) program package, invoking gradient geometry optimization (Schlegel, 1982). The optimized structural parameters were used in the vibrational frequencies calculations at the DFT/B3LYP/6-31G (d,p) level to characterize all stationary points as minima. The natural bonding orbitals (NBO) calculation (Glendering et al., 1998) wear performed using NBO 3.1 program as implemented in Gaussian 03W (Frisch et al., 2004) package at DFT/6-31G(d,p) level in order to understand various second-order interaction between the filled orbitals of one subsystem and vacant of another subsystem, which is measure of the intermolecular delocalization or hyper conjugation. The Raman activities (Si) calculated with Gaussian 03were converted to relative Raman intensity using Raint program (Michalska, 2003) by the expression:

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Fig.1. The atom numbering for TA molecule



Fig. 2. FTIR spectra of TA calculated and experimental

$$I_i = 10^{-12} (v_0 - v_i)^4 (1/v_i) S_i$$

Where t_i is the relative Raman intensity, S_i the Raman activities, v_i is the wave number of normal modes and v_0 denotes the wave number of the excitation laser (Michalska, 2005).

RESULTS AND DISCUSSION

Molecular geometry

The labeling of atoms in TA is given in Fig 3. The optimized geometrical parameters (bond length and bond angles) by DFT/B3LYP with 6-31G (d,p) basis set are listed in Table 1, a general priority for reproducing the experimental bond length



Fig. 3. FT-Raman spectra of TA calculated and experimental

taken from the Ref (Carl kemnitz, 2002) is not present among DFT/B3LYP levels. However, all the bond and bond angles computed with the DFT/B3LYP level shows excellent agreement with available experimental results.

Vibrational assignments

Vibrational spectral assignments have been performed on the recorded FT-IR and FT-Raman spectra based on the theoretically predicted wave numbers by density functional method (B3LYP) using 6-31G (d,p) basis set and have been collected in Table 2. Comparison of the frequencies calculated at B3LYP method with experimental values reveals the overestimation of the calculated vibrational modes due to neglect of anharmonicity in real system. In our study we have followed scaling factor for B3LYP/6-31G (d,p); 0.9608 (Scoott *et al.*, 1996).

C-H Vibration

The hetero aromatic structure shows the presence of C-H stretching vibration in the region 3100-3000 cm⁻¹ which is the characteristic region for the ready identification of C-H stretching vibration (Varsanyi, 1973). In this region, the bands are not affected appreciably by the nature of substituents. The vibration assigned by chemcraft program package at 2971 to 2952 cm⁻¹ by B3LYP/6-31G (d,p) levels, show good agreement with weak FT-Raman band at 2952 cm⁻¹ wear assigned to the C-H stretching vibration in the aromatic ring. In general in-plane and out-of-plane aromatic C-H deformations occur in the region 1300-1000 cm⁻¹ and 600-1000 cm⁻¹ respectively. The C-H in-plane bending vibration computed at 1057, 1196 and 1277 cm⁻¹ by B3LYP/6-31G(d,p) method shows good agreement with experimental values. The C-H out of plane bending vibration band appears in FT-IR at 845 cm⁻¹ and in FT-Raman at 852 cm⁻¹ shows good agreement with computed DFT (B3LYP) method at 838 cm⁻¹. In general C-H vibration computed by both methods shows good agreement with experimental observation as well as literature data.



Fig. 4. The atomic orbital HOMO – LUMO composition of the frontier molecular orbital for TA

environment, so they show pronounced shifts in the spectra of the hydrogen bonded species. The hydroxyl stretching vibrations are generally (Sajan *et al.*, 2006) observed in the region around 3500 cm⁻¹. In the case of the un-substituted phenols it has been shown that the frequency of O-H stretching vibration in the gas phase is 3657 cm^{-1} (Michalska *et al.*, 1996). Similarly in our case a strong FT-IR band at 3423 cm⁻¹ is assigned to O-H stretching vibration. The hydrogen bonding effect through hydroxyl group leads to dimmer conformation OH stretching mode calculated at 3652 cm^{-1} which is much closer to the FT-IR experimental observation at 3423 cm^{-1} .

The O-H in-plane bending vibration in the phenols, in general lies in the region 1150-1250 cm⁻¹ and is not much affected due to hydrogen bonding unlike to stretching and outof-plane bending frequencies. The medium band in FT-IR spectrum at 1162 and 1195 cm⁻¹ is assigned O-H in-plane bending vibration. Theoretically computed value at 1143 and 1163 cm⁻¹ by B3LYP method shows good agreement with recorded spectrum. The O-H out-of-plane bending mode for the free molecule lies below 300 cm⁻¹ and it is beyond the infrared spectral range of the present investigation. However, for the associated molecule (Varsanyi, 1973) the O-H out-of plane bending mode lies in the region 517-710 cm⁻¹ in both intermolecular and intramolecular association, the frequency is at a high value than in free O-H. In our present investigation a strong band observed in FT-IR spectrum at 555 cm⁻¹ is

 Table 1. Comparison between the calculated DFT and experimental values of geometrical parameters for TA

Bond length	B3LYP/6-31G(d,p) experimental	Bond angle	B3LYP/6- 31G(d,p)	experimental
C1-C2	1.545	1.523	C1-C7-O9	116.2	119.9
C1-C6	1.545	1.523	C3-C2-H13	110.4	109.4
C1-C7	1.526	1.509	C3-C2-H14	108.9	109.5
C1-H12	1.095	1.113	C2-C3-C4	112.2	109.4
C2-C3	1.536	1.523	C2-C3-H15	109.6	109.4
C2-H13	1.096	1.113	C2-C3-H16	109	109.4
C2-H14	1.1	1.113	H13-C2-H14	106.3	105.3
C3-C4	1.538	1.523	C4-C3-H15	110.3	109.4
C3-H15	1.097	1.113	C4-C3-H16	109.1	109.5
C3-H16	1.1	1.113	C3-C4-C5	110.5	109.5
C4-C5	1.538	1.523	C3-C4-C10	111.4	109.5
C4-C10	1.535	1.523	C3-C4-H17	107.5	109.5
C4-H17	1.104	1.113	H15-C3-H16	106.4	105.3
C5-C6	1.534	1.523	C5-C4-C10	112.1	109.4
C5-H18	1.093	1.113	C5-C4-H17	107.8	109.5
C5-H19	1.1	1.113	C4-C5-C6	112.1	109.5
C6-H20	1.096	1.1129	C4-5-H18	109.6	109.5
C6-H21	1.101	1.113	C4-C5-H19	108.8	109.5
C7-08	1.206	1.208	C10-C4-H17	107.4	109.5
C7-09	1.365	1.338	C4-C10-N11	111.3	109.5
O9-H22	0.968	0 9719	C4-C10-H23	108.8	109.4
C10-N11	1 467	1 4379	C4-C10-H24	108.7	109.5
C10-H23	1 098	1 113	C6-C5-H18	110.8	109.5
C10-H24	1 105	1 113	C6-C5-H19	109	109.5
N11-H25	1 017	1.02	C5-C6-H20	110.5	109.5
N11-H26	1 018	1.02	C5-C6-H21	108.9	109.5
			H18-C5-H19	106.4	105.3
Bond angle	B3LYP/6- 31G(d.p)	experimental	H20-C6-H21	106.3	105.3
C2-C1-C6	110.5	109.5	O8-C7-O9	119.6	119.9
C2-C1-C7	112.1	109.4	C7-O9-H22	110.1	109.5
C2-C1-H12	108.4	109.5	N11-C10-H23	107.7	109.4
C1-C2-C3	111.3	109.5	N11-C10-H24	113.8	109.4
C1-C2-H13	109.9	109.4	C10-N11-H25	110.1	109.4
C1-C2-H14	109.8	109.5	C10-N11-H26	109.7	109.4
C6-C1-C7	112.7	109.5	H23-C10-H24	106.3	105.3
C6-C1-H12	108.4	109.5	H25-N11-H26	106.1	109.4
C1-C6-C5	111.2	109.5	1120 1111-1120	100.1	107.4
C1-C6-H20	109.9	109.5			
C1-C6-H21	109.9	109.5			
C7-C1-H12	104.6	109.5			
C1-C7-08	124.3	119.9			
0.07-00	147.J	117.7			

O-H vibrations

The O-H group gives rise to three vibration (stretching, inplane bending and out-of-plane bending vibration). The O-H group vibrations are likely to be most sensitive to the assigned to O-H out-of-plan bending vibration, the theoretically computed value by B3LYP shows the same kind of vibration at 520 and 646 cm⁻¹ are assigned to O-H out-of-plane bending vibration.

Table 2. Vibrational wavenumbers obtained for TA at B3LYP/6-31G(d,p) and compared with experimental value

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$\begin{array}{cccccccccccccccccccccccccccccccccccc$	DH I ₂ I ₂ I ₂
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	I ₂ I ₂ I ₂
217 236 17 1 7-C-NH 253 276 1 3 7-C-NH	I ₂ I ₂ I ₂
253 276 1 3 γ -C-NF	I2 I2 I2
	I ₂
330 324 8 2 y-C-NH	
381 402 6 1 βOH	
420 434 3 1 γ CCC	2
452 442 16 3 γ CCC	2
472 475 467 71 1 γ CCC	
526 496 5 3 $\gamma OH+\rho C$	H_2
555 520 519 8 0 ρ-CH ₂	2
008 043 / 0 γOH	
769 750 8 13 BCC	
788 767 772 2 2 BCC	
790 799 59 4 v CH	
845 845 837 57 3 y CH	
875 21 1 v CC	
881 6 0 υ CH	
903 911 6 2 υCC	
921 920 926 4 2 γ CH	
955 964 8 2 β CCC	2
988 8 6 βCC	
1010 1008 1004 3 9 βCH	
1029 1052 2 Ι βCH	
1059 1058 1056 1 2 BCH	
1050 1050 1 2 p CH 1070 2 6 B CH	
1120 1120 1119 1 6 B CH	
1141 1 3 в Он	
1162 1154 1161 4 2 β ОН	
1195 1196 1195 6 1 β CH	
1227 1232 1221 14 13 β CH ₂	
1255 1254 1243 12 19 β CH ₂	
$1263 6 / \beta CH_2$	2 T
1270 335 4 UC-OF	1
1280 1270 1275 52 2 pCH	
1200 1 peri 1296 4 2 t CH ₂	
1302 1302 2 2 t CH	
1330 1325 1329 0 4 β OH	
1339 0 4 v CC	
1344 2 6 vCO	
1356 1351 1357 1 3 8 CH ₂	
1385 $13/4$ $13/9$ 15 3 $1CH_2$ 1434 1436 0 21 BCH	
1441 4 1 BCH	
1445 1446 14 1 B CH	
1453 1455 0 8 t CH ₂	
1538 1465 1463 2 17 δ NH ₂ +υ	CC
1635 1652 1642 25 12 υ C=O+υ	CC
1799 294 14 υ C=C)
2842 2830 83 85 $v_s CH_2$	2
2804 2808 2855 10 $/1$ $U_{s}CH_{2}$	2
2070 2000 2005 10 54 $U_8 CH_2$ 2801 37 146 5 CH	2
$2901 2902 9 21 110 0_{s} CH_{2}$	2
2907 53 38 D.CH	<u>.</u>
2921 2932 2932 35 53 U ₈ CH	2
2949 52 64 v CH	
2952 2952 5 92 v CH	
2959 33 69 v CH	
2967 43 88 v CH	
2989 17 87 v CH	
3352 I 140 U ₈ NH 3423 2425 0 105 U ₈ NH	ſ
3648 18 45 DO	L

υ- stretching; υ_s- symmetric stretching; υ_{as}- asymmetric stretching; β- in-plane bending; γ- out of plane bending; δ-scissoring; ω-wagging; ρ- rocking; t- twisting; τ- tor

Table 3. Second order perturbation theory analysis of Fork matrix in NBO basis for TA

Donor(i)	Туре	ED/e	Acceptor(j)	Туре	ED/e	E(2) (kcal/mol)
C4 C10		1.00	C2 C2	*	015	2.16
C4-C10	σ	1.96	02-03	σ.	.015	2.16
			C3-C4	σ^*	.019	0.99
			C4-C5	σ^*	0.19	1.07
			C4-H17	σ^*	0.02	0.55
			N11-H25	σ^*	0.006	2.13
N11-H25	σ	1.98	C4-C10	σ^*	0.02	3.09
N11-H26	σ	1.98	C10-H23	σ^*	0.018	3.03
LP(1)N11		1.96	C4-C10	σ^*	0.02	0.87
			C10-H23	σ*	0.01	1.12
			C10-H24	σ^*	0.03	7.39
LP(1)O9		1.82	C1-C7	σ^*	0.06	7.46
			C7-O8	σ^*	0.01	2.24
LP(2)O9		1.82	C7-C8	σ*	0.20	47.40
LP(1)O8		1.97	C1-C7	σ*	0.06	2.61
			C7-O9	σ*	0.09	1.01
LP(2)O8		1.855	C1-C7	σ*	0.022	17.07
			C7-O9	σ*	0.09	34.57

Table 4. HOMO LUMO energy calculated by B3LYP/6-31G(d,p) method

Method	B3LYP/6-31G(d,p)
HOMO	-11.323eV
LUMO	0.369eV
Energy gap (ΔE)	-11.692eV

C-C ring stretching

In benzonitriles, the distance between two carbon atoms changes the ring angles because of its substituents groups such .as cyanogens. There are six equivalent C-C bonds in benzene and consequently there will be six C-C stretching vibrations. The bands are observed at 1538 and 1632 cm⁻¹ in FT-IR are identified as C-C stretching vibrations. The same vibration appear in the FT-Raman spectrum at 1652 cm⁻¹. The theoretically scaled C-C stretching vibrations by DFT method are at 1642⁻¹ shows excellent agreement with recorded FT-IR and FT-Raman spectral data. The ring in-plane vibrations have given rise to weak bands across the low frequency region, that is to say, below 1000cm⁻¹ the bands at 751 and 989 cm⁻¹ have been assigned to C-C in-plane bending vibrations.

CH₂ bending vibrations

In the present study, the band around 1280cm⁻¹ in FT-IR is assigned to CH₂ wagging mode which agrees with the result (1281cm⁻¹) of (Matulkova et al) in the case of 4-aminotriazole adipic acid (4atadip). The harmonic frequencies calculated by B3LYP/6-31G(d,p) method falling in the region 1340-1277 cm⁻¹, which are in agreement with the earlier report (Matulkova et al). The calculated frequencies 1297 and 1359 cm^{-1} are assigned to twisting of CH_2 . Matulkova et al. reported the CH₂ twisting frequency observed in the range 1223-1314 cm-1 of 4atadip. In the present work CH₂ twisting vibration is assigned to medium Raman band at 1276 and 1302 cm⁻¹. However, the DFT value is in agreement with the experimental value. The band appeared around 1356 cm-1 in FT-IR is assigned to CH2 scissoring mode. The computed CH2 rocking mode is appears at 520 cm⁻¹ in DFT. The FT-IR band at 526 cm⁻¹ and Raman band at 535 cm⁻¹ are assigned to CH₂ rocking mode, which agree favourably with Matulkova et at. and also find support from theoretical value. Billes et al. assigned CH-in-plane bending vibration appeared at 1419, 1114, 1411 and 1126 cm⁻¹, respectively in 1H-1,2,3-and 1D-1,2,3-triazoles. In view of above bands appeared at 1435 cm⁻¹

in FT-Raman and 1445 in FT-IR spectrum are assigned to C-H in plane bending mode. The calculated frequencies in the region 828-956 cm⁻¹ for C-H out-of- plane bending fall in the FT-IR and FT-Raman values of 845-920 cm⁻¹.

NH₂ vibration

The molecule under investigation possesses only one NH₂ group and hence one expects one symmetric and one asymmetric N-H stretching vibration in NH₂ group. In all the primary aromatic amines, the N-H stretching frequency occurs in the region 3300-3500 cm⁻¹ (Bellamy, 1980). Hence the weak bands in FT-IR spectrum wear located at 3320-3232 cm⁻ assingned to N-H asymmetric and symmetric stretching vibration, respectively in NH2 group. These assignments agree well with the earlier report (Bellamy, 1980). The scaled NH₂ asymmetric and symmetric stretching are in the range 3439-3356 cm⁻¹ in B3LYP/6-31G(d,p). The computed NH₂ scissoring vibration at 1465 cm⁻¹ in B3LYP/6-31G(d,p) is in agreement with the expected experimental value at 1465 cm⁻¹. The C-NH₂ out-of-plane and in-plane bending vibration at 253 and 381 cm⁻¹ observed at FT-Raman spectrum agree well with theoretically obtained values using B3LYP/6-31G (d,p).

COOH vibration

Carboxylic acid dimmer is formed by strong hydrogen bonding in the solid and liquid state vibrational analysis of carboxylic acid group is made on the basis of carbonyl group and hydroxyl group. The C=O stretch of carboxylic acid is identical to the C=O stretch in ketones, which is expected in the region 1740-1660 cm⁻¹ (Vein et al., 1991). The C=O bond formed by P_{π} - P_{π} between C and O intermolecular hydrogen bonding reduces the frequencies of the C=O stretching absorption to a greater degree than does intermolecular H bonding because of the different electro-negativities of C and O, the bonding are not equally distributed between the two atoms. The loan pair of electrons on oxygen also determines the nature of the carbonyl groups. In our present study a strong band observed in FT-IR spectrum 1635 cm⁻¹ is assigned to C=O stretching vibrations, which show good agreement with B3LYP scaled value at 1642 cm⁻¹. Two other characteristic carboxylic group vibrations are C-O stretching and C-O bending vibrations. They are expected in the region 1140-395 cm⁻¹ and 1700-875 cm⁻¹ depending on whether monomeric, dimeric or other hydrogen bonded species are present. Generally the C-O stretching mode appears at lower frequencies than C-O bending vibration.

NBO Analysis

Natural bond orbital analysis gives the most accurate possible natural Lewis structure picture of \emptyset because all orbital are mathematically chosen to include the highest possible percentage of the electron density. Interaction between both filled and virtual orbital space information is correctly explain by the NBO analysis, it could enhance the analysis of intra and inter molecular interaction. The second order Fock matrix was carried out to evaluate donor (i)-acceptor (j) i.e. donar level bonds to acceptor level bond interaction in the NBO analysis (Szafran *et al.*, 2007). The result of interaction is a loss of occupancy from the concentration of electron NBO of the idealized Lewis structure into an empty non-Lewis orbital. For each donor (i) and acceptor (j), the stabilization energy E(2) associates with the delocalization $i \rightarrow j$ is estimated as

$$E^{(2)} = \Delta E_{ij} = q_i (F(i,j)^2 / (\varepsilon_j - \varepsilon_i))$$

Where qi is the donor orbital occupancy, ε_i - and ε_i are diagonal element and F(i,j) is the off diagonal NBO Fock matrix element. Natural bond orbital analyses provide an efficient method for interaction among bond, and also provide a convenient basis for investigation charge transfer or conjugative interaction in molecular systems. Some electron donor orbital, acceptor orbital and the interaction stabilization energy resulted from the second order micro-disturbance theory are reported (James et al., 2006 and Jun-na et al., 2005). The larger $E^{(2)}$ value the more intensive is the interaction between electron donors and acceptor i.e. the more donation tendency from electron donors to electron acceptors and the greater the extent of conjugation of the whole system (Sebastian et al., 2010). Delocalization of electron density between occupied Lewis-type (bond or lone pair) NBO orbitals and formally unoccupied (anti bond or rydberg) non Lewis NBO orbital correspond to astabilizing donor-acceptor interaction. NBO analysis has been performed on the Ta molecule at the DFT/B3LYP/6-31G(d,p) level in order to elucidate, the intra molecular rehybridization and delocalization of electron density within the molecule. The LP(2) O9 is seen to be lowest-occupancy and to be primarily delocalized into ant bond C_7 - O_8 . The $E^{(2)}$ values and types of transition as shown Table 3.

HOMO and LUMO analysis

Highest occupied molecular orbital (HOMO) and the lowest unoccupied molecular orbital (LUMO) and very important parameters for quantum chemistry. We can determine the way the molecule interacts with other species. Hence, they are called the frontier orbitals. HOMO, which can be thought the outermost orbital containing electrons, trends to give these electrons such as an electron donor. On the other hand; LUMO can be thought the innermost orbital containing free places to accept electrons (Gece et al., 2008). The HOMO and LUMO energy calculated by B3LYP /6-31G(d,p) method are shown in Table 4. This electronic transition absorption corresponds to the transition from the ground to the first excited state and is mainly described by an electron excitation from the HOMO to the LUMO. The HOMO is located over the NH₂ group, the HOMO > LUMO transition implies an electron density transfer to ring and acid group. The atomic compositions of the frontier molecular orbital are shown in Fig. 4. The calculated self-consistent field (SCF) energy of TA is -519.11 A.U

Conclusion

The spectral studies such as FT-IR, FT-Raman for TA was carried out for the first time. A complete vibrational and molecular structure analysis has been performed based on the quantum mechanical approach by DFT (B3LYP) calculations. The assignments made at DFT level of theory with only reasonably deviations from the experimental values seem to be correct. The calculated absorption maxima values at DFT (B3LYP) level almost correlate with the experimental value. The theoretically constructed FT-IR, FT-Raman spectra

coincide with experimentally observed FT-IR, FT-Raman. NBO reflects the charge transfer within the molecule. HOMO and LUMO orbitals have been visualized.

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