## RESEARCH ARTICLE

# STEADY FLOWS IN PIPES OF RECTANGULAR CROSS-SECTION THROUGH POROUS MEDIUM WITH MAGNETIC FIELD 

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## ARTICLE INFO

## Article History:

Received $28^{\text {th }}$ December, 2014
Received in revised form
$15^{\text {th }}$ December, 2014
Accepted $27^{\text {th }}$ January, 2015
Published online $26^{\text {th }}$ February, 2015

## Key words:

Steady flow,
Rectangular cross section, Incompressible fluid, porous medium and Magnetic field.


#### Abstract

In this paper we have investigated the steady flow in pipes of rectangular cross-section through porous medium with magnetic field. We have investigated the velocity, flux and vortex line.


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## INTRODUCTION

We have investigated the steady flow in pipes of rectangular cross-section through porous medium. Attempts have been made by several researchers i.e. Givler and Altobelli (1994) a determination of the effective viscosity for the Brinkman-Forchheimer flow model. Goel and Agarwal (1998) shears flow instability of visco- elastic fluid in a porous medium. Goldstein (1930) concerning some solutions of the boundary-layer equations in hydrodynamics.

Gopinath (1994) steady streaming due to small amplitude superposed oscillations of a sphere in a viscous fluid. Gorski and Bernard (1995) Vorticity Transport Analysis of turbulent flows. Journal of Fluids Engng. Goyeau and Gobin (1996) Numerical study of double-diffusive natural convection in a porous cavity using the Darcy-Brinkman formulation. Goyon (1996) highReynolds number solutions of Navier-stokes equations using incremental unknowns. Grigoriev and Dargush (1999) a polyregion boundary element method for incompressible viscous fluid flows.

Grosan, Revnic, Pop and Ingham (2009) Magnetic field and internal heat generation effects on the free convection in a rectangular cavity filled with a porous medium. Gupta and Manohar (1979) Boundary approximations and accuracy in viscous flow computations. Gupta (1991) high accuracy solutions of incompressible Navier-stokes equations. Haddad and Corke (1998) Boundary layer receptivity to free-stream sound on parabolic bodies. Haddon and Riley (1985) on flows with closed streamlines. Haji-Sheikh and Vafai (2004) analysis of flow and heat transfer in porous media imbedded inside various-shaped ducts. In this paper we have investigated the velocity, flux and vortex line.

## FORMULATION OF THE PROBLEM

Let z -axis be taken the direction of flow along the axis of the pipe. Then $\mathrm{u}=0, \mathrm{v}=0$ for steady and incompressible fluid the velocity component is independent of $\mathbf{z}$.

[^0]The equation of continuity. $\quad \frac{\partial u}{\partial x}+\frac{\partial v}{\partial y}+\frac{\partial w}{\partial z}=0$ $\qquad$
But $u=0, v=0, \frac{\partial w}{\partial z}=0 \Rightarrow w=w(x, y) \ldots \ldots$.
i.e. $w$ is independent of $\mathbf{z}$

The Navier-Stokes equations of motion in the absence of body forces.
$-\frac{\partial P}{\partial x}=0$
$-\frac{\partial P}{\partial y}=0$.
$\frac{\partial P}{\partial z}+\mu\left(\frac{\partial^{2} w}{\partial x^{2}}+\frac{\partial^{2} w}{\partial y^{2}}\right)-\left(\frac{1}{\rho K}+\frac{\sigma B_{0}^{2}}{\rho \mu}\right) \mu w=0$.


Figur-1

It is clear from (3) and (4) $\mathbf{P}$ is independent of $\mathbf{x}$ and $\mathbf{y}$ i.e. $\mathbf{p}$ is the Function of $\mathbf{z}$
SOLUTION OF THE PROBLEM $\quad \mathrm{p}=\mathrm{p}(\mathrm{z}) \frac{\partial p}{\partial z}=\frac{d p}{d z}=$ Constant $=-\mathrm{P}$
let $\left(\frac{1}{\rho K}+\frac{\sigma B_{0}^{2}}{\rho \mu}\right)=B^{2}, \mu\left[\frac{\partial^{2} w}{\partial x^{2}}+\frac{\partial^{2} w}{\partial y^{2}}-B^{2} w\right]=\frac{d p}{d z} \Rightarrow \frac{\partial^{2} w}{\partial x^{2}}+\frac{\partial^{2} w}{\partial y^{2}}-B^{2} w=-\frac{P}{\mu} \ldots \ldots \ldots \ldots$.
$\left(D^{2}+D^{\prime 2}-B^{2}\right) w=-\frac{P}{\mu} \quad \therefore \quad C . F=\sum a_{n} e^{h_{n} x+h_{n}^{\prime} y}$ where $h_{n} \& h_{n}^{\prime}$ are related by $h_{n}{ }^{2}+h_{n}{ }_{n}{ }^{2}-B^{2}=0$
and P.I. $=\frac{1}{D^{2}+D^{\prime 2}-B^{2}}\left(-\frac{P}{\mu}\right)=\frac{P}{B^{2} \mu} \Rightarrow \mathrm{w}(\mathrm{x}, \mathrm{y})=\sum_{n=1}^{\infty} a_{n} e^{h_{n} x+h_{n}^{\prime} y}+\frac{1}{B^{2} \mu} P$ Where $h_{n}^{2}+h_{n}^{2}=B^{2}$

Case-1 $\quad w(x, y)=0$ at $(a, b) w(x, y)=0$ at $(a,-b)$
$\sum_{n=1}^{\infty} a_{n} e^{h_{n}{ }^{a}+h_{n}^{\prime} b}+\frac{P}{\mu B^{2}}=0 \quad$ and $\quad \sum_{n=1}^{\infty} a_{n} e^{h_{n}{ }^{a}-h_{n}^{\prime} b}+\frac{P}{\mu B^{2}}=0$
$\Rightarrow-\frac{P}{\mu B^{2}}=\sum_{n=1}^{\infty} a_{n} e^{h_{n}{ }^{a}+h_{n}^{\prime} b} \ldots \ldots . .(a) \&-\frac{P}{\mu B^{2}}=\sum_{n=1}^{\infty} a_{n} e^{h_{n}{ }^{a}-h_{n}^{\prime} b}$.
on solving $h_{n}^{\prime}=0 \Rightarrow h_{n}=-B \Rightarrow-\frac{P}{\mu B^{2}}=e^{-a B} \sum_{n=1}^{\infty} a_{n} \Rightarrow \sum_{n=1}^{\infty} a_{n}=e^{-\frac{P}{\mu B^{2}}} e^{-a B}$
$w_{1}(x, y)=-\frac{P}{\mu B^{2}} e^{a B} e^{-x B}+\frac{P}{\mu B^{2}} \quad \Rightarrow-\frac{P}{\mu B^{2}}=e^{-a B} \sum_{n=1}^{\infty} a_{n} \quad \Rightarrow \sum_{n=1}^{\infty} a_{n}=e^{-\frac{P}{\mu B^{2}}} e^{-a B}$
Case-2 $\quad w(x, y)=0 \quad$ at $(-a, b) \&(-a,-b)$
$w_{2}(x, y)=-\frac{P}{\mu B^{2}} e^{a B} e^{x B}+\frac{P}{\mu B^{2}}=-\frac{P}{\mu B^{2}} e^{B(x+a)}+\frac{P}{\mu B^{2}}$

## Case - 3

$w(x, y)=0$ at $(-a, b) \&(a, b) \Rightarrow w_{3}(x, y)=-\frac{P}{\mu B^{2}} e^{b B} e^{-y B}+\frac{P}{\mu B^{2}}=-\frac{P}{\mu B^{2}} e^{B(-y+b)}+\frac{P}{\mu B^{2}}$
Case-4 $\quad w(x, y)=0$ at $(-a,-b) \quad \&(a,-b) \Rightarrow w_{4}(x, y)=-\frac{P}{\mu B^{2}} e^{B(y+b)}+\frac{P}{\mu B^{2}}$
$w(x, y)=\frac{P}{\mu B^{2}}\left[1-2 e^{a B} \operatorname{Cosh} x B-2 e^{b B} \operatorname{Cosh} y B\right]$.
In particular case: In the case of square i.e $. a=b$
$w(x, y)=\frac{P}{\mu B^{2}}\left[1-2 e^{a B}(\operatorname{Cosh} x B+\operatorname{Cosh} y B)\right]$.

## Flux $\mathbf{Q}$ of the fluid over an area of rectangular cross-section:

$Q=\int_{x=-a}^{a} \int_{y=-b}^{b} w(x, y) d x d y=\int_{-a}^{a} \int_{-b}^{b} \frac{P}{\mu B^{2}}\left\{1-2 e^{a B} \operatorname{Cosh} x B-2 e^{b B} \operatorname{Cosh} y B\right\} d x d y$
$=\frac{2 P}{\mu B^{2}} \int_{-a}^{a} \int_{0}^{b}\left\{1-2 e^{a B} \operatorname{Cosh} x B-2 e^{b B} \operatorname{Cosh} y B\right\} d y d x=\frac{2 P}{\mu B^{2}} \int_{-a}^{a}\left\{\left(1-2 e^{a B} \operatorname{Cosh} x B\right) b-2 e^{b B}\left(\frac{1}{B} \operatorname{Sinh} b B\right)\right\} d x$ $=\frac{4 P}{\mu B^{2}} \int_{0}^{a}\left\{b\left(1-2 e^{a B} \operatorname{Cosh} x B\right)-\frac{2}{B} e^{b B} \operatorname{Sinh} b B\right\} d x=\frac{4 P}{\mu B^{2}}\left[b\left\{x-\frac{2}{B} e^{a B} \operatorname{Sinh} x B\right\}_{0}^{a}-\frac{2 a}{B} e^{b B} \operatorname{Sinh} b B\right]$
$=\frac{4 P}{\mu B^{2}}\left[b\left\{a-\frac{2}{B} e^{a B} \operatorname{Sinh} a B\right\}-\frac{2 a}{B} e^{b B} \operatorname{Sinh} b B\right]$
$Q=\frac{4 P}{\mu B^{2}}\left[a b-\frac{2}{B}\left(b e^{a B} \operatorname{Sinh} a B+a e^{b B} \operatorname{Sinh} b B\right)\right]$

In particular case: In the case of square $a=b$
$Q=\frac{4 P}{\mu B^{2}}\left[a^{2}-\frac{4 a}{B} e^{a B} \operatorname{Sinh} a B\right]$

The equation of vortex line: $\frac{d x}{\Omega_{x}}=\frac{d y}{\Omega_{y}}=\frac{d z}{\Omega_{z}}$ here $\Omega_{x}, \Omega_{y} \& \Omega_{z}$ are vorticity components
where $\bar{q}=u i+v j+w k=\frac{P}{\mu B^{2}}\left[1-2 e^{a B} \operatorname{Cosh} x B-2 e^{b B} \operatorname{Cosh} y B\right] \hat{k}$
$\Omega_{x}=\frac{\partial w}{\partial y}-\frac{\partial v}{\partial z}=\frac{P}{\mu B^{2}}\left[-2 B e^{b B} \operatorname{Sinh} y B\right]=-\frac{2 P}{\mu B} e^{b B} \operatorname{Sinh} y B$
$\Omega_{y}=\frac{\partial u}{\partial z}-\frac{\partial w}{\partial x}=-\frac{P}{\mu B^{2}}\left[-2 B e^{a B} \operatorname{Sinh} x B\right]=\frac{2 P}{\mu B} e^{a B} \operatorname{Sinh} x B \quad \Omega_{z}=0$
$\Rightarrow \frac{d x}{-\frac{2 P}{\mu B} e^{b B} \operatorname{Sinh} y B}=\frac{d y}{\frac{2 P}{\mu B} e^{a B} \operatorname{Sinh} x B}=\frac{d z}{0} \Rightarrow d z=0 \Rightarrow z=B$
$\frac{d x}{-e^{b B} \operatorname{Sinh} y B}=\frac{d y}{e^{a B} \operatorname{Sinh} x B} \Rightarrow e^{a B} \int \operatorname{Sinh} x B d x+e^{b B} \int \operatorname{Sinh} y B d y=C_{1}$
$\frac{1}{B} e^{a B} \operatorname{Cosh} x B+\frac{1}{B} e^{b B} \operatorname{Cosh} y B=C_{1}$ or $e^{a B} \operatorname{Cosh} x B+e^{b B} \operatorname{Cosh} y B=C_{1} B=A$
$\therefore$ Vortex lines: $\quad e^{a B} \operatorname{Cosh} x B+e^{b B} \operatorname{Cosh} y B=A \quad \& \quad \mathrm{Z}=\mathrm{B}$

## Clearly the flow is Rotational in pipe.

I
n particular case: In the case of square $a=b$

$$
\begin{equation*}
e^{a B}[\operatorname{Cosh} x B+\operatorname{Cosh} y B]=A \quad \& \quad \mathrm{Z}=\mathrm{B} \tag{12}
\end{equation*}
$$

Tables for velocity: Case- 1
Let $P=\frac{1}{4}, \mu=.5, \quad a=b=1, \frac{1}{\sqrt{\sigma K}}=\sqrt{\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{1}{2} \Rightarrow B=\sqrt{\frac{1}{\sigma K}+\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{1}{\sqrt{2}}$

Table 1 (for velocity)

|  | $(x, y)$ | $(.1, .1)$ | $(.2, .3)$ | $(.3, .4)$ | $(.4, .5)$ | $(.5, .6)$ | $(.6, .7)$ | $(.7, .8)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\frac{1}{\sqrt{\rho K}}=\frac{1}{2}$ | $w(x, y)$ | -11.21 | -11.297 | -11.396 | -11.529 | -11.696 | -11.897 | -12.133 |
| $\sqrt{\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{1}{2}$ | $w(x, y)$ | -11.21 | -11.297 | -11.396 | -11.529 | -11.696 | -11.897 | -12.133 |
| $\sqrt{\frac{1}{\rho K}+\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{1}{\sqrt{2}}$ | $w(x, y)$ | -7.133 | -7.245 | -7.367 | -7.532 | -7.739 | -7.99 | -8.286 |

Case- 2
Let $P=\frac{1}{4}, \quad \mu=.5, \quad a=b=1, \frac{1}{\sqrt{\sigma K}}<\sqrt{\frac{\sigma B_{0}^{2}}{\rho \mu}}$ Let $\frac{1}{\sqrt{\sigma K}}=\frac{1}{2}$ and $\sqrt{\frac{\sigma B_{0}^{2}}{\rho \mu}}=1 \Rightarrow B=\sqrt{\frac{1}{\sigma K}+\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{\sqrt{5}}{2}$
Table- 2 (for velocity)

|  | $(x, y)$ | $(.1, .1)$ | $(.2, .3)$ | $(.3, .4)$ | $(.4, .5)$ | $(.5, .6)$ | $(.6, .7)$ | $(.7, .8)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\frac{1}{\sqrt{\rho K}}=\frac{1}{2}$ | $w(x, y)$ | -11.21 | -11.297 | -11.396 | -11.529 | -11.696 | -11.897 | -12.133 |
| $\sqrt{\frac{\sigma B_{0}^{2}}{\rho \mu}}=1$ | $w(x, y)$ | -4.964 | -5.114 | -5.28 | -5.504 | -5.788 | -6.134 | -6.547 |
| $\sqrt{\frac{1}{\rho K}+\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{\sqrt{5}}{2}$ | $w(x, y)$ | -4.525 | -4.695 | -4.882 | -5.135 | -5.458 | -5.854 | -6.328 |

Case - 3
Let $P=\frac{1}{4}, \mu=.5, \quad a=b=1, \frac{1}{\sqrt{\sigma K}}>\sqrt{\frac{\sigma B_{0}^{2}}{\rho \mu}}$ Let $\frac{1}{\sqrt{\sigma K}}=1$ and $\sqrt{\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{1}{2} \Rightarrow B=\sqrt{\frac{1}{\sigma K}+\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{\sqrt{5}}{2}$
Table- 3 (for velocity)

|  | $(x, y)$ | $(.1, .1)$ | $(.2, .3)$ | $(.3, .4)$ | $(.4, .5)$ | $(.5, .6)$ | $(.6, .7)$ | $(.7, .8)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\sqrt{\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{1}{2}$ | $w(x, y)$ | -4.964 | -5.114 | -5.28 | -5.504 | -5.788 | -6.134 | -6.547 |
| $\frac{1}{\sqrt{\rho K}}=1$ | $w(x, y)$ | -11.21 | -11.297 | -11.396 | -11.529 | -11.696 | -11.897 | -12.133 |
| $\sqrt{\frac{1}{\rho K}+\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{\sqrt{5}}{2}$ | $w(x, y)$ | -4.525 | -4.695 | -4.882 | -5.135 | -5.458 | -5.854 | -6.328 |

## CONCLUSION AND DISCUSSION

In this paper we have investigated the velocity by the Table-1 of equations (7) between velocities and point $(x, y)$ it is clear that paths of velocity are approximately parallel increases uniformly with negative sign in porous medium, magnetic field and porous medium with magnetic field in the interval $(.1, .1) \leq(x, y) \leq(.7, .8)$ but the value of velocity in porous medium and magnetic field at $\frac{1}{\sqrt{\rho K}}=\sqrt{\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{1}{2}$ is greater than the corresponding value of velocity in porous medium with magnetic field at $\sqrt{\frac{1}{\rho K}+\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{1}{\sqrt{2}}$ in the interval $(.1, .1) \leq(x, y) \leq(.7, .8)$.
Again by the Table-2 of equations (7) between velocities and point $(x, y)$ it is clear that paths of velocity are approximately parallel increases uniformly with negative sign in porous medium, magnetic field and porous medium with magnetic field in the interval $(.1, .1) \leq(x, y) \leq(.7, .8)$ but the value of velocity in porous medium at $\frac{1}{\sqrt{\rho K}}=\frac{1}{2}$ is greater (negatively) then corresponding value of velocity in magnetic field at $\sqrt{\frac{\sigma B_{0}^{2}}{\rho \mu}}=1$ and also is greater than the corresponding value of velocity in porous medium with magnetic field at $\sqrt{\frac{1}{\rho K}+\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{\sqrt{5}}{2}$ respectively.
Again by the Table-3 of equations (7) between velocities and point $(x, y)$ it is clear that paths of velocity are approximately parallel increases uniformly with negative sign in porous medium, magnetic field and porous medium with magnetic field in the interval $(.1, .1) \leq(x, y) \leq(.7, .8)$ but the value of velocity in magnetic field at $\sqrt{\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{1}{2}$ is greater (negatively) then corresponding value of velocity in porous medium at $\frac{1}{\sqrt{\rho K}}=1$ and also is greater than the corresponding value of velocity in porous medium with magnetic field at $\sqrt{\frac{1}{\rho K}+\frac{\sigma B_{0}^{2}}{\rho \mu}}=\frac{\sqrt{5}}{2}$ respectively. We have investigated the vortex lines and the volumetric flow of elliptic and circle given by the equations (8), (9), (10), (11) and (12) respectively.

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